

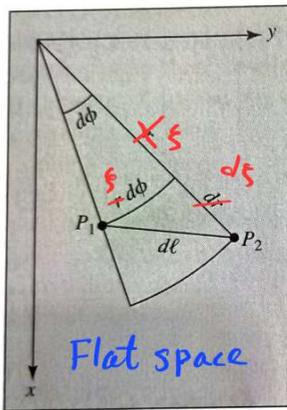
☆ Spatial Part of Metric

$$ds^2 = -c^2 dt^2 + a^2(t) d\mathbf{x}^2$$

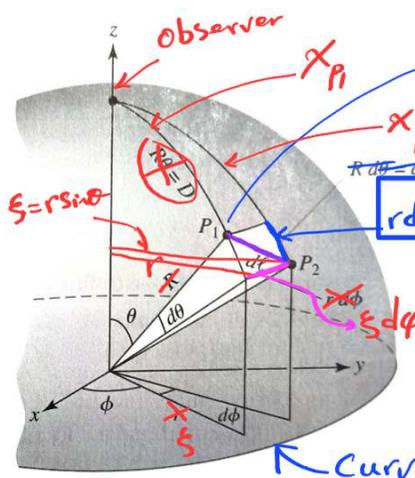
my goal is to derive the general form of $g_{\mu\nu}$ for static and 3D space Homogeneous Universe.

① First approach: 2D space (spherical surface embedded in 3D)

Chapter 29 Cosmology



(a)



(b)

FIGURE 29.18 dl as measured for (a) a flat plane and (b) the surface of a sphere.

$P_1(x_1, y_1, z_1), |r_{P_1}| = \sqrt{x_1^2 + y_1^2 + z_1^2} = ct_1$
 $P_2(x_2, y_2, z_2), |r_{P_2}| = \sqrt{x_2^2 + y_2^2 + z_2^2} = ct_2$
 $x^2 + y^2 + z^2 = r^2$
 $r d\theta = dx$
 $d\ell_{2D}^2 = ?$
 Due to 2D

My Notation $R \equiv$ Physical Distance (proper distance) scale Factor
 $\chi \equiv$ Comoving Distance $R(t) = a(t) \chi$

$\chi = r\theta$, $\xi = r \sin\theta$

$d\ell_{2D}^2 = |r_{P_1} - r_{P_2}|^2 = r^2 d\theta^2 + \xi^2 d\phi^2$

$\xi = r \sin\theta \rightarrow \frac{d\xi}{r \sin\theta} = r d\theta = dx$

$d\ell_{2D}^2 = \frac{ds^2}{(1 - \frac{\xi^2}{r^2})} + \xi^2 d\phi^2$

$\sin\theta = \frac{\xi}{r}$

$\cos^2\theta = 1 - \sin^2\theta = 1 - \frac{\xi^2}{r^2}$

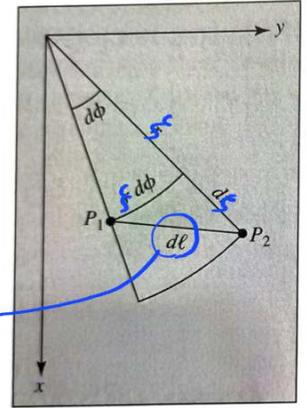
$K \equiv \frac{1}{r^2}$: Curvature
 انحناء

For 2D space $dl_{2D}^2 = \frac{d\xi^2}{1-K\xi^2} + \xi^2 d\varphi^2$

$$dl_{2D}^2 = dx^2 + S^2(x) d\varphi^2$$

$x \rightarrow \xi$ for $K=0$ Flat space.

$$dl_{2D}^2 = d\xi^2 + \xi^2 d\varphi^2$$



Now 3D space (Spherical space embedded in 4D) $dl_{3D}^2 = ?$

② ☆ Second approach 2D: $dl_{2D}^2 = ?$

$$dl^2 = dx^2 + dy^2 + dz^2$$

$$r^2 = x^2 + y^2 + z^2 \rightarrow \text{2D space, } r = \text{fixed}$$

$$0 = x dx + y dy + z dz$$

$$z dz = -(x dx + y dy) \quad \text{sphere}$$

$$dz = -\frac{x dx + y dy}{z} = \pm \frac{x dx + y dy}{\sqrt{r^2 - x^2 - y^2}} \quad \text{Hypersphere}$$

$$x^2 + y^2 = \xi^2 \rightarrow \begin{cases} x dx + y dy = \xi d\xi \rightarrow (x dx + y dy)^2 = \xi^2 d\xi^2 \\ dx^2 + dy^2 = d\xi^2 + \xi^2 d\varphi^2 \end{cases}$$

$$dl_{2D}^2 = dx^2 + dy^2 + dz^2 = dx^2 + dy^2 + \frac{(x dx + y dy)^2}{r^2 - (x^2 + y^2)}$$

$$= d\xi^2 + \xi^2 d\varphi^2 + \frac{\xi^2 d\xi^2}{r^2 - \xi^2}$$

$$= d\xi^2 \left[1 + \frac{\xi^2}{r^2 - \xi^2} \right] + \xi^2 d\varphi^2 = d\xi^2 \left[\frac{r^2 - \xi^2 + \xi^2}{r^2 - \xi^2} \right] + \xi^2 d\varphi^2$$

$$= \frac{d\xi^2}{1 - \frac{\xi^2}{r^2}} + \xi^2 d\varphi^2 = \frac{d\xi^2}{1 - k\xi^2} + \xi^2 d\varphi^2$$

$$dl_{2D}^2 = dx^2 + S^2(x) d\varphi^2$$

For 3D : $dl_{3D}^2 = ?$

$$dl_{2D}^2 = d\xi^2 + \xi^2 d\varphi^2 + \frac{\xi^2 d\xi^2}{r^2 - \xi^2}$$

$r^2 = x^2 + z^2 = \text{const}$

$$dl_{3D}^2 = d\xi^2 + \xi^2 (d\theta^2 + \sin^2\theta d\varphi^2) + \frac{\xi^2 d\xi^2}{r^2 - \xi^2}$$

$$dl_{3D}^2 = \frac{d\xi^2}{1 - k\xi^2} + \xi^2 d\Omega^2$$

$$dl_{3D}^2 = dx^2 + S^2(x) d\Omega^2$$

$$ds^2 = -c^2 dt^2 + a^2(t) \left[dx^2 + S^2(x) d\Omega^2 \right]$$

or $dl_{4D}^2 = dx^2 + dz^2$ $x^2 + z^2 = r^2$ شعاع کره چینه بندی

4D → 3D, r = fixed

$$dl_{3D}^2 = dx^2 + K \frac{(x \cdot dx)^2}{1 - Kx^2}$$

(Recall $dl_{2D}^2 = dx^2 + \frac{r^2 d\theta^2}{r^2(1 - \frac{\theta^2}{r^2})}$)

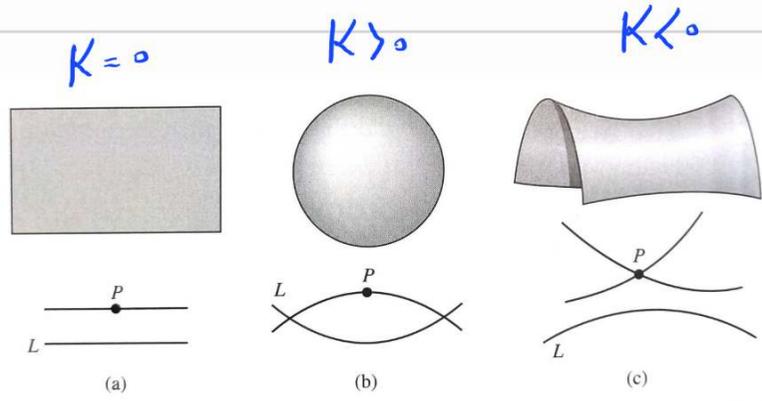
$$ds^2 = -cdt^2 + a^2(t) dl_{3D}^2$$

What about $S(x) = ?$

$$\left. \begin{aligned} dx &= \frac{d\xi}{1 - K\xi^2} \\ \int dx &= \int \frac{d\xi}{1 - K\xi^2} \end{aligned} \right\} \rightarrow x = \begin{cases} \xi & K=0 \\ \frac{1}{\sqrt{|K|}} \sinh^{-1}(\sqrt{|K|}\xi) & K < 0 \\ \frac{1}{\sqrt{|K|}} \sin^{-1}(\sqrt{|K|}\xi) & K > 0 \end{cases}$$

$$\rightarrow \xi = \begin{cases} x & K=0 \\ \frac{1}{\sqrt{|K|}} \sinh(\sqrt{|K|x}) & K < 0 \\ \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|x}) & K > 0 \end{cases}$$

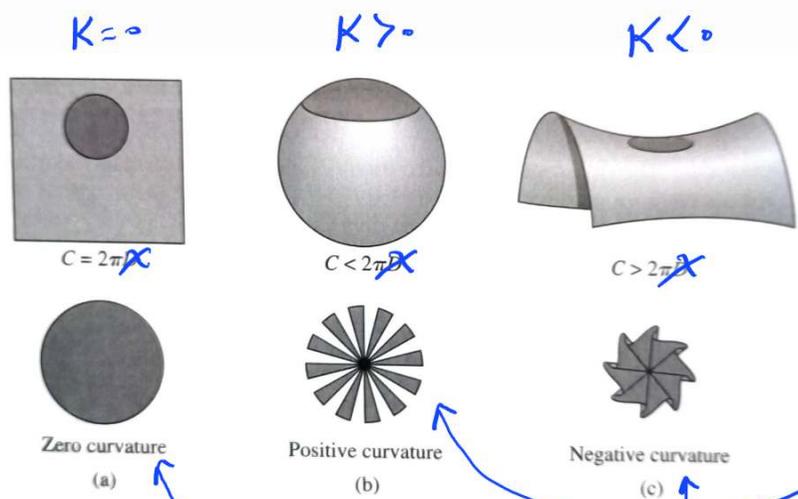
$$\xi = S(x) = \frac{1}{\sqrt{|K|}} \text{Sinn}(\sqrt{|K|x})$$



برای فرضیه اقلیدس
در فضای تخت از نقطه P
تنها یک خط موازی با آن
در فضای منبسطه (K>0) هیچ خط
موازی وجود ندارد. در فضای

FIGURE 29.15 The parallel postulate, illustrated for three alternative geometries: (a) Euclidean, (b) elliptic, and (c) hyperbolic.

منبسطه (K<0) بی نهایت خط وجود دارند موازی با آن.



$$C = \int S(x) d\varphi$$

$$= 2\pi S(x)$$

$$C < 2\pi x$$

$$C = 2\pi x$$

$$C > 2\pi x$$

FIGURE 29.17 Calculating the curvature of a surface in three geometries: (a) a flat plane, (b) the surface of a sphere, and (c) the surface of a hyperboloid.

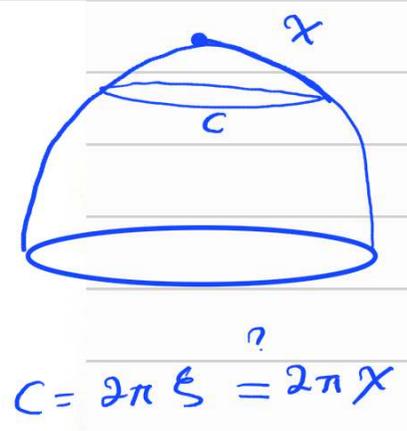
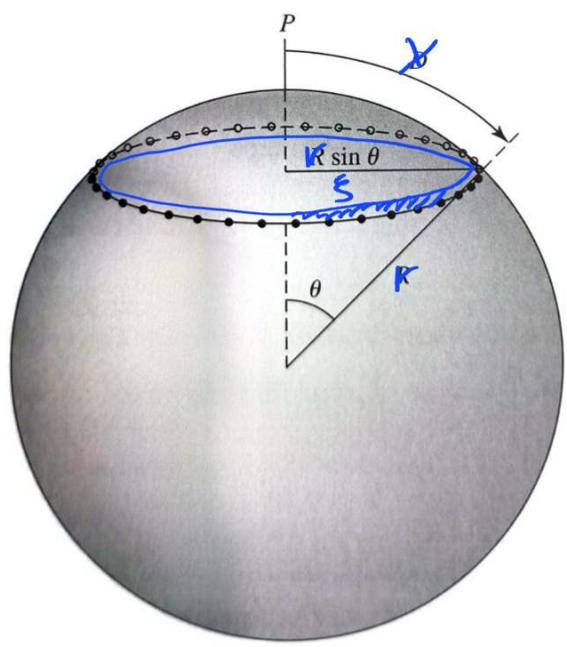


FIGURE 29.16 A local measurement of the curvature of a sphere.

$$dl_{2D}^2 = dx^2 + S^2(x) d\varphi$$

$$dx=0 \rightarrow dl_{2D} = dC = S(x) d\varphi$$

↑
Circumference

$$C = \int_0^{2\pi} d\varphi S(x) = S(x) 2\pi$$

$$K=0 \quad S(x)=x \rightarrow C_{K=0} = 2\pi x = 2\pi r$$

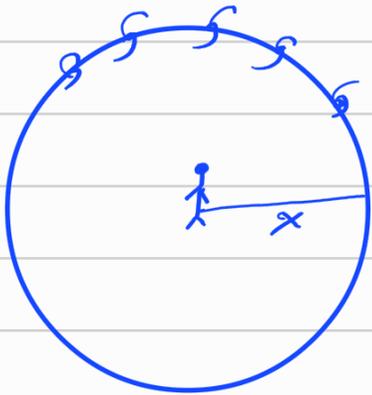
$$K>0 \quad C_{K>0} = \int_0^{2\pi} \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|} x) d\varphi = 2\pi < 2\pi x$$

$$K<0 \quad C_{K<0} = \int_0^{2\pi} \frac{1}{\sqrt{|K|}} \sinh(\sqrt{|K|} x) d\varphi = 2\pi > 2\pi x$$

$$dA_{2D} = dl_{2D}^2 = dx^2 + S(x)^2 d\varphi^2$$

$$dA_{2D} = C_{2D}(x) dx$$

$$= \int_0^x dx' C_{2D}(x') = \pi x^2 \quad \text{For } K=0$$



کبریا را نشان ندهی بزین تعیین انحنای بیان به سطح زیرین

در صورتی که فاصله بین ستاره که این است در آن صورت

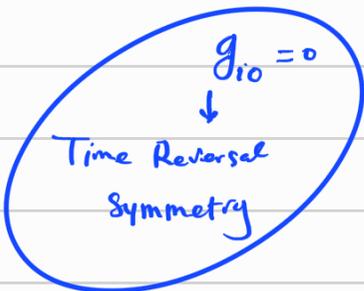
$$N_g < N_g < N_g$$

↑ ↑ ↑
K>0 K=0 K<0

closed ← Flat ← open

Exercise 1:

For static and homogeneous space and using Schwarzschild metric show that



$$dl^2 = a^2(t) [f(r) dr^2 + r^2 d\Omega^2]$$

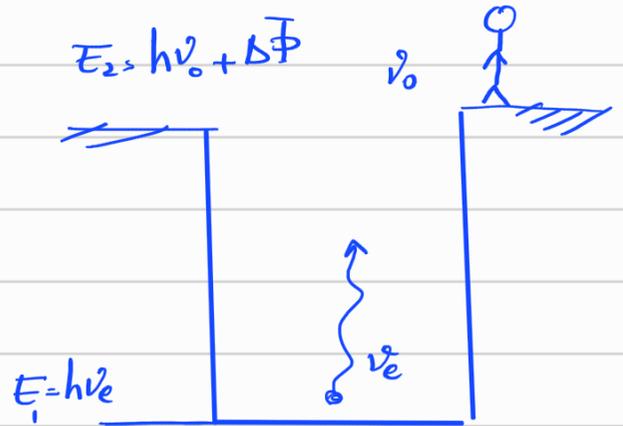
$$f(r) = \frac{1}{1-Kr^2}$$

Ricci scalar should be constant

$$R = g^{\mu\nu} R_{\mu\nu} = c t s$$

☆ Exercise 2: Gravitational Redshift

$$\nu_e \xrightarrow{\circ} \nu_o$$

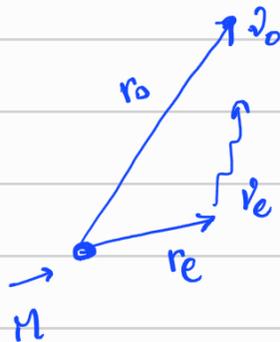


$$E_2 = E_1 - \Delta\phi \quad \text{از دریا پایین} \quad \& \quad \Delta\phi = \phi_2 - \phi_1$$

$$h\nu_o = h\nu_e \left(1 - \frac{\Delta V}{c^2}\right)$$

$$\Delta V \approx \frac{\Delta\phi}{m}$$

$$\frac{\nu_o}{\nu_e} = 1 - \frac{\Delta V}{c^2} \Rightarrow 1+z = 1 + \frac{GM}{c^2} \left(\frac{1}{r_e} - \frac{1}{r_o} \right)$$

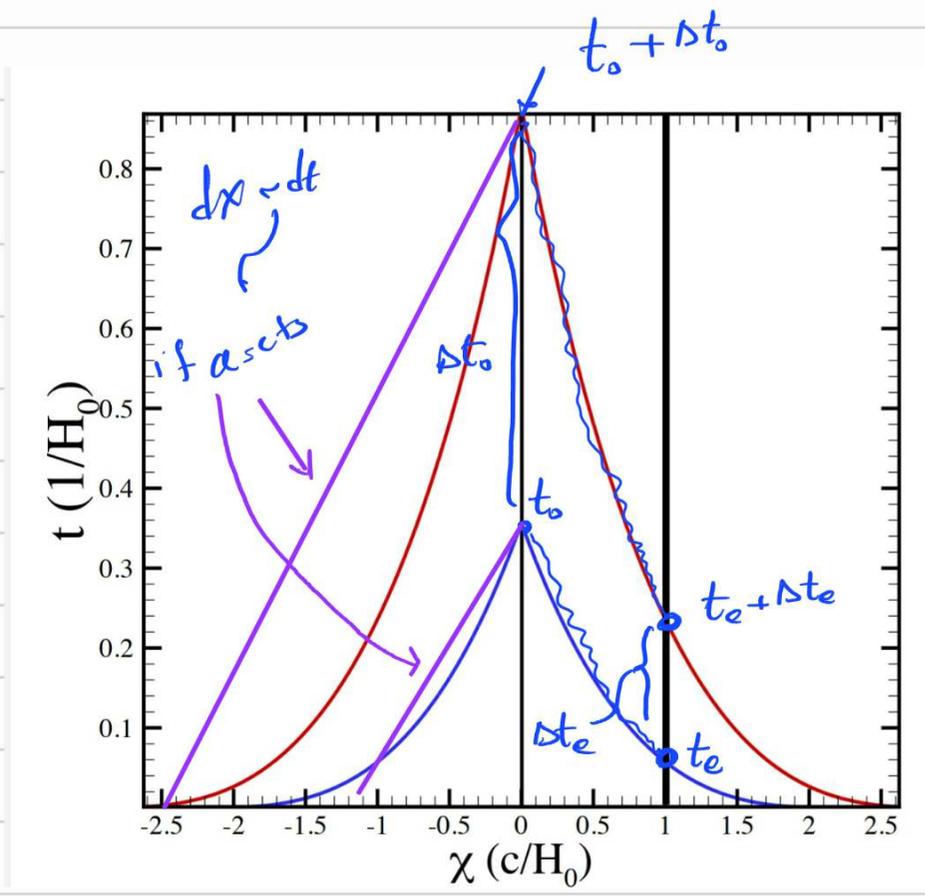


$$\begin{aligned} r_e &\gg \frac{2MG}{c^2} \\ r_o &\gg \frac{2MG}{c^2} \end{aligned}$$

☆ Cosmological Redshift

$$ds^2 = -c^2 dt^2 + a^2(t) \left[dx^2 + S^2(x) d\Omega^2 \right]$$

$$\left(\begin{array}{l} \text{Null geodesic} \\ \& \\ d\Omega = 0 \end{array} \right) \rightarrow ds^2 = 0 \rightarrow \boxed{dx = \frac{dt}{c a(t)}}$$



$$\chi = c \int_{t_e}^{t_o} \frac{dt}{a(t)} = c \int_{t_e + \Delta t_e}^{t_o + \Delta t_o} \frac{dt}{a(t)}$$

because the relative location of observer and source does not change. $\chi = \text{fixed}$

$$\int_{t_e}^{t_o} = \int_{t_o + \Delta t_o}^{t_e + \Delta t_e} + \int_{t_e}^{t_o} + \int_{t_o}^{t_o + \Delta t_o}$$

$$c \int_{t_o}^{t_o + \Delta t_o} \frac{dt}{a(t)} = c \int_{t_e}^{t_e + \Delta t_e} \frac{dt}{a(t)}$$

$$\frac{\Delta t_o}{a(t_o)} = \frac{\Delta t_e}{a(t_e)}$$

$$\Delta t_o = \frac{a(t_o)}{a(t_e)} \Delta t_e$$

Suppose that
 $a(t_o) = 1$

$$v_o = a(t_o) v_e = \frac{1}{1+z} v_e$$

$$\lambda_o = (1+z) \lambda_e$$

$$\frac{\Delta \lambda}{\lambda} = z$$

- Physical length
- Comoving length
- Angular diameter distance
- Luminosity distance
- Comoving volume element
- Cosmic age

$$R(t) = a(t)\chi$$

$$ds^2 = c^2 dt^2 - a(t)^2 [d\chi^2 + S(\chi) d\Omega^2]$$

$$\begin{cases} ds^2 = 0 \\ d\Omega^2 = 0 \end{cases} \rightarrow d\chi = \frac{cdt}{a(t)} \rightarrow \chi = c \int_t^{t_0} \frac{dt'}{a(t')} = c \int_0^z \frac{dz}{H(z)}$$

$$\begin{cases} ds^2 = 0 \\ d\chi^2 = 0 \end{cases} \rightarrow a(t)S(\chi)\theta = \Delta l \rightarrow d_A \equiv a(t)S(\chi)$$

$$d_l \equiv (1+z)S(\chi) = d_A(1+z)^2$$

$$f \equiv \frac{dV}{d\Omega dz} = \frac{S(\chi)^2}{H(z)} \rightarrow \frac{dN}{dz} = n(z) \frac{dV}{dz} = 4\pi n(z) \frac{S(\chi)^2}{H(z)}$$

$$\Delta N_{L > L_{\min}} = \int_{L_{\min}}^{\infty} \frac{dN}{dz} dl \Delta z$$

$$t_0 = \int_0^{t_0} dt = \int_0^{\infty} \frac{dz}{(1+z)H(z)}$$

Now, I turn to compute some quantities from Background Evolution which are related to Space-Time distance (ds)

① Comoving Distance to a Source at z



$$z_{obs} = 0$$

$$a_{obs} = 1$$

○ Horizon $dx = \frac{c dt}{a(t)}$ $\chi_0 = \int_0^{\chi_0} dx = \int_0^{t_0} c \frac{dt}{a(t)}$

$$\chi_0 = c \int_0^{\infty} \frac{dz}{H(z)}$$

$$\begin{cases} z_0 = 0 \rightarrow a_0 = 1 \\ z_{\text{Big-Bang}} = \infty \rightarrow a_{\text{Big-Bang}} = 0 \end{cases}$$

$$\chi(z) = c \int_0^z \frac{dz'}{H(z')} = \frac{c}{H_0} \int_0^z \frac{dz'}{H(z')}$$

$$H_0 = h \cdot 100 \frac{\text{Km}}{\text{s}} / \text{Mpc}$$

$$\chi(z) = \frac{3 \times 10^5 \text{ Km/s} \cdot \text{Mpc}}{100 h \text{ Km/s}} \int_0^z \frac{dz'}{H(z')}$$

$$\chi(z) = 3000 \left(\frac{\text{Mpc}}{h} \right) \int_0^z \frac{dz'}{H(z')}$$

$$\chi_0 = c \int \frac{dt}{a(t)} \quad \text{if } a = ct_0 \quad \chi_0 \approx ct_0$$

$$a(t_1) < a(t_2) < a(t_3) \dots < a(t_0)$$

$$\chi_0 > ct_0$$

چونکہ تیز رفتاری سے دور سے دور کے لیے 14×10^{10} yr ، اس لیے کہ 13.8×10^9 yr عام ہے

○ Event Horizon $d_E = \int_{t_0}^{\infty} \frac{c dt}{a(t)} = \begin{cases} \infty & \text{No event Horizon} \\ \neq \infty & \text{Event Horizon} \end{cases}$

در حالتی که زمان مورد نیاز برای

درافت اطلاعات از یک جسم که در فاصله d_E باشد بی نهایت شود. یعنی جسم دور نمی آید.

Ex: For de Sitter Universe $\Lambda > 0$ and $H \propto \sqrt{\Lambda}$

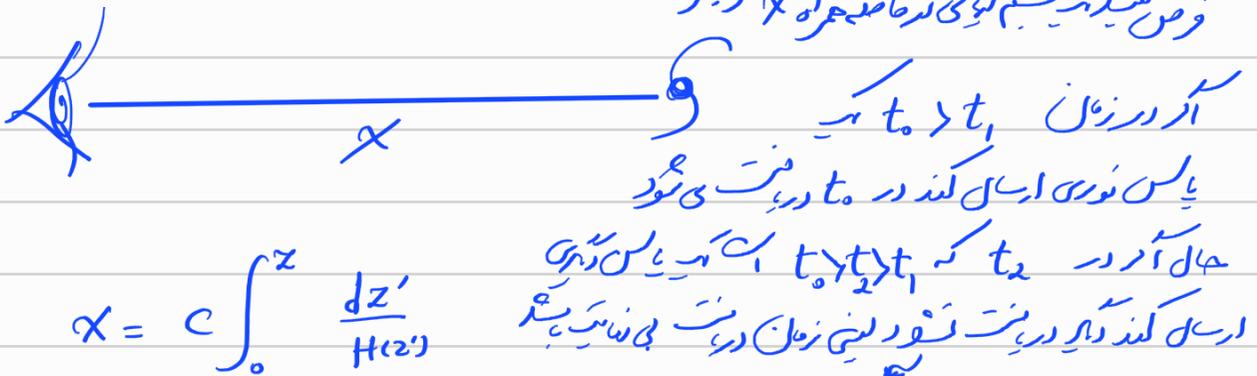
$\underline{a(t) = e^{H_0 t}}$ $\dot{a} = H_0 e^{H_0 t}$ $\frac{\dot{a}}{a} = H_0$ ← **Cosmological Inflation era**

$$d_E = \int_t^{\infty} \frac{c dt'}{e^{H_0 t'}} = \int_t^{\infty} c dt' e^{-H_0 t'} = \frac{c}{H_0} \neq \infty$$

★ Maximum Visible age of Cosmological object.

بیشترین سنی که یک جسم کیهانی را می بینیم

فرض کنید یک جسم کیهانی در فاصله x قرار دارد



$$x = c \int_0^z \frac{dz'}{H(z')}$$

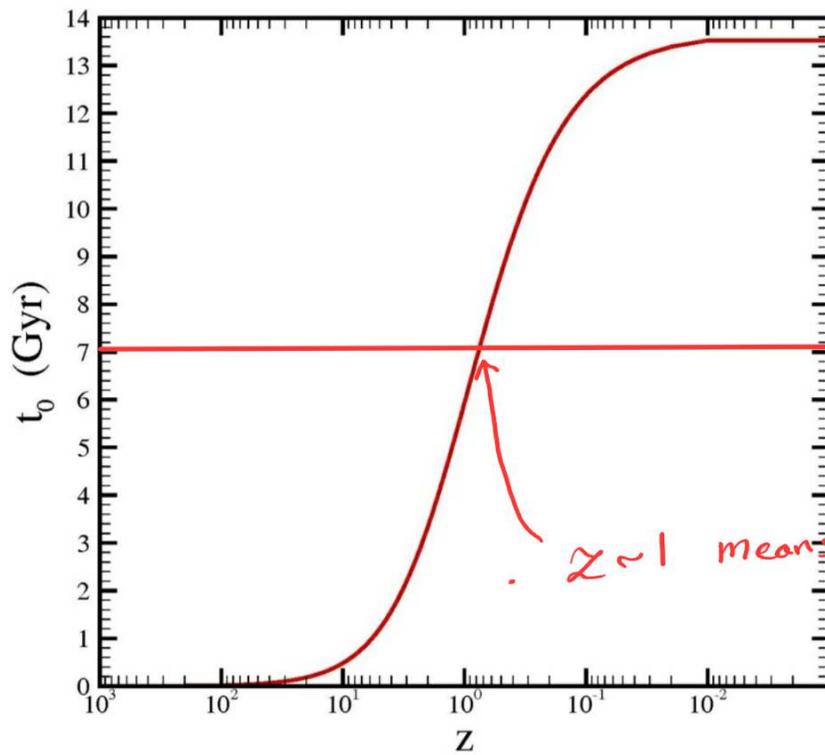
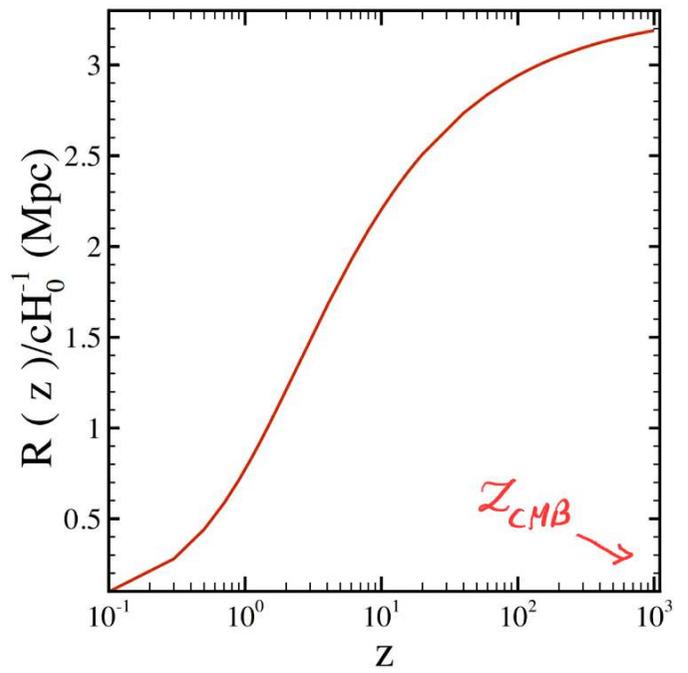
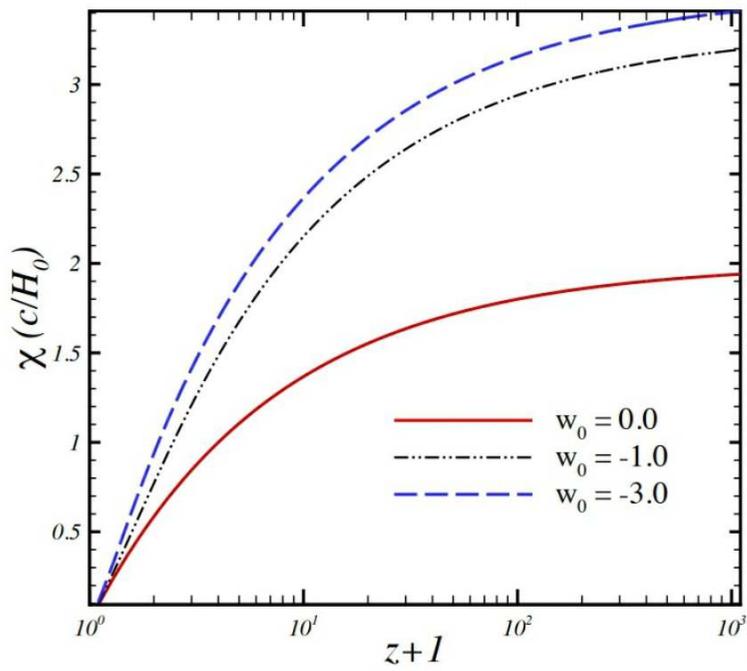
$$= c \int_{t_1}^{t_0} \frac{dt}{a(t)} = \int_{t_2}^{t_0 + \Delta t_0} \frac{dt}{a(t)} = \int_{t_{max}}^{\infty} \frac{dt}{a(t)}$$

$$t_0 > t_2 > t_1$$

$$T_{hidden} = t_{max} - t_1$$

↓
 This means that if we can see an object which is located at x , after passing time

about $T_{hidden}(x)$, we are not able to see it.



$$t \approx \frac{t_0}{2}$$

Maximum visible age of astronomical objects

PRD, 65,
047301
(2002)

$$\int_{t_1}^{t_0} \frac{dt}{a(t)} \cong \int_{t_2}^{\infty} \frac{dt}{a(t)} \rightarrow t_2 = t_{\max} \rightarrow t_{\max} - t_1 = T_{\text{hidden}}$$

