

Review on Perturbative RG (Goldenfeld)

Chapter 12

① $Z = \int Ds e^{\mathcal{H}}$ while $Z = \int Dm e^{-\beta \mathcal{H}}$

$\mathcal{H} = \int d^d r \left[\frac{1}{2} (\nabla S)^2 + \frac{1}{2} r_0 S^2 + \frac{1}{4} u_0 S^4 - h_0 S \right]$ (12.5), (12.61)

Kardar $\beta \mathcal{H} = \beta \mathcal{H}_0 + \mathcal{U} = \int d^d r \left[\frac{K}{2} (\nabla m)^2 + \frac{t}{2} m^2 + \frac{1}{2} (\nabla^2 m)^2 + \dots + \mathcal{U}(m \cdot m)^2 \right]$
 ↳ (5.1) Kardar.

F.T. $-\tilde{\mathcal{H}} = \frac{1}{2} \int \frac{d^d K}{(2\pi)^d} (r_0 + K^2) |S(K)|^2$ (12.65)

$+ \frac{1}{4} u_0 \int \frac{d^d k_1 d^d k_2 d^d k_3 d^d k_4}{(2\pi)^{4d}} (2\pi)^d \delta_D(k_1 + k_2 + k_3 + k_4)$

$\times S(k_1) S(k_2) S(k_3) S(k_4)$

$-\tilde{\mathcal{H}} = -\mathcal{H}_0 - V(\{S\}, \{\sigma\}) = -\mathcal{H}_0(\{S\}) - \mathcal{H}_0(\{\sigma\}) - V$

F.T. $\tilde{\beta \mathcal{H}} = \int \frac{d^d q}{(2\pi)^d} \frac{t + Kq^2 + Lq^4 + \dots}{2} |m(q)|^2$

$+ u \int \frac{d^d q_1 d^d q_2 d^d q_3 d^d q_4}{(2\pi)^{4d}} (2\pi)^d \delta_D(q_1 + q_2 + q_3 + q_4)$

$m_\alpha(q_1) m_\alpha(q_2) m_\beta(q_3) m_\beta(q_4)$

$\alpha, \beta = 1, \dots, n$
 n -vector field

② Green's function $\langle S(k_1) S(k_2) \rangle_0 = \frac{(2\pi)^d \delta_D(k_1 + k_2)}{r_0 + k_1^2}$ (12.96)

$$\langle m_\alpha(\varphi) m_\beta(\varphi') \rangle_0 = \frac{(2\pi)^d \delta_D(\varphi + \varphi') \delta_{\alpha\beta}}{t + K\varphi^2 + L\varphi^{4+}} \quad (5.6) \quad \text{Kardar}$$

③ RG: Coarse-graining and rescaling and Renormalizing.

$$\begin{array}{ccc}
 & & \begin{array}{l} K' = lK \\ q' = l^d q \end{array} \\
 & \downarrow & \begin{array}{l} S' = \frac{S}{Z} \\ \tilde{m}' = \tilde{m}/Z \end{array} \\
 S \equiv \{ S'(K) \} \oplus \{ \sigma(K) \} & & m \equiv \{ \tilde{m}(q) \} \oplus \{ \sigma(q) \} \\
 & \downarrow \quad \downarrow & \text{Kardar} \\
 \circ \langle K \rangle \sim \gamma_e & \gamma_e \langle K \rangle \sim 1 & \\
 \underbrace{a \langle r \rangle \sim \infty}_{\text{Singular Part}} & \underbrace{a \langle r \rangle \sim la}_{\text{Regular Part (Integrate out)}} &
 \end{array}$$

$$\textcircled{4} \quad Z = Z_0 \int Ds' e^{\mathcal{H}_s'} \langle e^{\tilde{V}} \rangle_0 \quad (12.74)$$

$$= Z_0 \int Ds' e^{\mathcal{H}_s'} e^{\langle \tilde{V} \rangle_0 + \frac{1}{2} [\langle \tilde{V}^2 \rangle_0 - \langle \tilde{V} \rangle_0^2] + \mathcal{O}(\tilde{V}^3)}$$

$\underbrace{\mathcal{O}(\tilde{V}^2)}_{\text{Combination of Coarse-graining + Perturbat}}$

$$(12.77) \quad \mathcal{H} = \mathcal{H}_\sigma + \mathcal{H}_s + \langle \tilde{V} \rangle_0 + \frac{1}{2} [\langle \tilde{V}^2 \rangle_0 - \langle \tilde{V} \rangle_0^2] + \mathcal{O}(\tilde{V}^3)$$

$$(5.32) \quad \tilde{\beta}\mathcal{H} = V \delta f_\ell^0 + \int_0^{M\ell} \frac{d^d q}{(2\pi)^d} \left(\frac{t + Kq^2}{2} |\tilde{m}(q)|^2 - \ln \langle e^{-u} \rangle_0 \right)$$

$$(5.33) \quad \ln \langle e^{-u} \rangle_0 = - \langle u \rangle_0 + \frac{1}{2} [\langle u^2 \rangle_0 - \langle u \rangle_0^2] + \mathcal{O}(u^3)$$

$$(12.80) \quad V \equiv \frac{1}{4} u_0 \int \frac{d^d k_1 d^d k_2 d^d k_3 d^d k_4}{(2\pi)^{4d}} (2\pi)^d \delta_D(k_1+k_2+k_3+k_4) S(k_1) S(k_2) S(k_3) S(k_4)$$

$\underbrace{\hspace{10em}}$
 $\left. \begin{array}{l} n=1 \text{ with respect} \\ \text{to Kadan} \end{array} \right\}$

⑤ $\langle V \rangle_0 = ?$ Section (12.3.4)

$u_0^{(1)} \rightarrow @ (V^2)$
 $u_0^{(2)} \rightarrow l$
 s^2

$$t' = t_l = z^2 l^{-d} \left[r_0 l^2 + K_l^2 + 2 \times 6 \underbrace{\text{Diagram}} \right] \quad (12.112)$$

Recall $\frac{1}{2} r_0 s^2$
 $s s s s \rightarrow s^2$

$u_0^{(4)} \rightarrow @ (V^3)$
 $u_0^{(4)} \rightarrow l$
 s^4

$$u_0' = z^4 l^{-3d} \left[\text{Diagram} \right] \quad (12.113)$$

$$\begin{aligned} \mathcal{Q}(V^2) \quad \eta_t = \chi_t = 2 \\ \eta_u = \chi_u = 4-d = \epsilon \end{aligned} \quad \left| \begin{array}{l} (t^*, u^*) = (0, 0) \\ \text{Gaussian fixed} \\ \text{Point} \end{array} \right.$$

For $d > 4 \rightarrow \eta_u < 0$ — Irrelevant
 compatible with perturbation

For $d < 4 \rightarrow \eta_u > 0$ — Relevant \rightarrow break down Perturbation

To mitigate this challenge
 we should compute up to $@(V^3)$

$$\frac{1}{2!} [\langle V^2 \rangle_0 - \langle V \rangle_0^2] = ?$$

$$? = I_1 \equiv \int_{N_e}^{\wedge} \frac{d^d k}{(2\pi)^d} \frac{1}{r_0 + k^2}$$

~~XX~~

$$I_2 \equiv \int_{N_e}^{\wedge} \frac{d^d k}{(2\pi)^d} \frac{1}{(r_0 + k^2)} \frac{1}{(r_0 + (k_1 + k_2 - k)^2)} \quad (12.130)$$

$k_1 = k_2 = 0$

$$? = I_2 \equiv \int_{N_e}^{\wedge} \frac{d^d k}{(2\pi)^d} \frac{1}{(r_0 + k^2)^2}$$

→ Long-mode

1974 Wilson

up to leading order near critical point

Symmetric Integrals

* Hint: $d^d q, d^d p \rightarrow d^d q$
(191, 192)

$$I = \int \frac{d^d q}{(2\pi)^d} f(q^2) = \frac{S_d}{(2\pi)^d} \int dq q^{d-1} f(q^2)$$

Surface area

$$S_d = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

$$\boxed{\epsilon = 4-d} \quad \text{---} \quad \boxed{d = 4-\epsilon}$$

Using (12.136) and (12.138) → (12.131) and (12.132)

⇓

$$r' = l^2 \left[r_0 + \frac{3}{4\pi^2} \frac{u_0}{4} \Lambda^2 \left(1 - \frac{1}{l^2} \right) - \frac{3}{2\pi^2} \frac{u_0}{4} r_0 \ln l \right] \quad (12.147)$$

$$\frac{u'}{4} = \frac{u_0}{4} + \frac{u_0}{4} \left[\epsilon - \frac{9}{2\pi^2} \frac{u_0}{4} \right] \ln l \quad (12.148)$$

$$l = (1 + \delta l) = e^{\delta l}$$

Fixed Points ?

Section (12.5.3)

$$l = 1 + \delta l \rightarrow \beta_l = 0 \Rightarrow r^*, u^* = ?$$

★ First Fixed Point $(r^*, u^*) = (0, 0)$ Gaussian fixed point

★ Second Fixed Point $(r^*, u^*) = \left(-\frac{\epsilon \Lambda^2}{6}, \frac{8\pi^2 \epsilon}{9} \right) \sim \mathcal{O}(\epsilon)$

We obtain that Second fixed point is related to Gaussian fixed point up to $\mathcal{O}(\epsilon)$

Wilson - Fisher Fixed Point

The Linearized RG transform

$$[K]_l = R_l[K] \rightarrow [K]_l = [K]_* + R \frac{\partial R}{\partial K} + \dots$$

$$K' = T R$$

$$T \leftarrow M v = \Lambda v$$

↙
eigen vector

↘
Eigen-value

$$M = \begin{pmatrix} \frac{\partial r'}{\partial r} & \frac{\partial r'}{\partial u} \\ \frac{\partial u'}{\partial r} & \frac{\partial u'}{\partial u} \end{pmatrix}$$

Fixed - pt { Gau
W.F.

$$(12.152) \quad M_{Gau} = \begin{pmatrix} l^2 & \frac{3\Lambda^2(l^2-1)}{4\pi^2} \\ 0 & l \end{pmatrix} \quad (12.152)$$

$$M_{W.F.} = \begin{pmatrix} l^{2-\frac{\epsilon}{3}} & \frac{3\Lambda^2(l^2-1)}{4\pi^2} \\ 0 & l \end{pmatrix} \quad (12.153)$$

$\Lambda = 1/a$

☆ At Gaussian Fixed Point

$(r^*, u^*) = (0, 0)$

Eigen-Value $\left\{ \begin{aligned} \Lambda_r = \Lambda_t = l^2 &\Rightarrow y_t = y_r = 2 \\ \Lambda_u = \Lambda_2 = l^d &\Rightarrow y_u = y_2 = d \end{aligned} \right.$

for $ISm d_c = 4$

$\epsilon > 0 \rightarrow d < 4$

Same Un. Verbal. by Class

$d > d_c \rightarrow$

Eigen Vector $\left\{ \begin{aligned} V_1 = (1, 0) &= (1r, 0 \times u) \\ V_2 = (-\frac{3\Lambda^2}{4\pi^2}, 1) \end{aligned} \right.$

(6.17) $d > \frac{2\beta + \gamma}{\nu} \equiv d_c$

Page 171 Goldenfeld

Upper Critical Dimension



According to Ginzburg Criterion at $d > 4$ is always satisfied - MF. ($\alpha < 8S^2$)

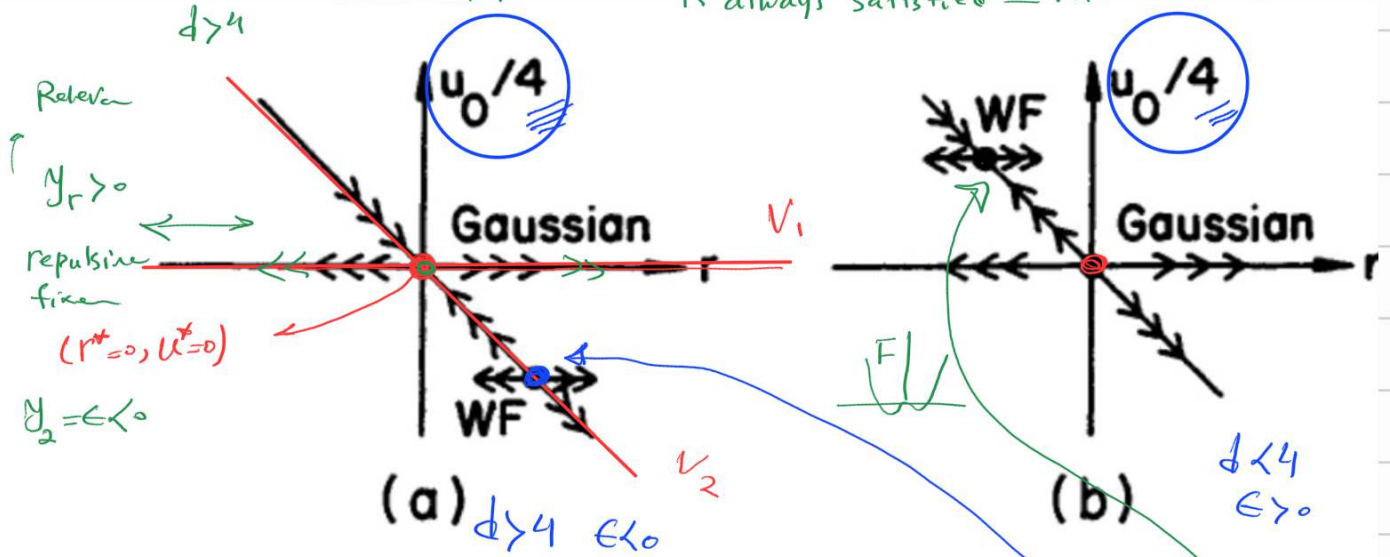


Figure 12.2 RG flows near four dimensions: (a) $d > 4$. (b) $d < 4$.

WF: $(r^*, u^*) = \left(-\frac{\epsilon \Lambda^2}{6}, \frac{8\pi^2}{9} \epsilon \right)$

$(r^*, \frac{u^*}{4}) = \left(-\frac{\epsilon \Lambda^2}{6}, \frac{2\pi^2}{9} \epsilon \right)$

For $d > 4$ ($\epsilon < 0$) $\rightarrow (r^*, \frac{u^*}{4}) = \left(+\frac{|\epsilon| \Lambda^2}{6}, -\frac{2\pi^2}{9} |\epsilon| \right)$

For $d < 4$ ($\epsilon > 0$) $\rightarrow (r^*, \frac{u^*}{4}) = \left(-\frac{|\epsilon| \Lambda^2}{6}, +\frac{2\pi^2}{9} |\epsilon| \right)$

For $d > 4$ ($\epsilon < 0$) $y_u < 0 \rightarrow$ Irrelevant

$$\text{W.F. } (r^*, u_{\frac{x}{4}}^*) = \left(\frac{+|\epsilon|\Lambda^2}{6}, -\frac{2\pi^2|\epsilon|}{9} \right)$$

$$u_{\frac{x}{4}}^* = -\frac{2\pi^2|\epsilon|}{9} < 0$$

Not acceptable



No physical solution

$$u^* < 0 \rightarrow M = \sqrt{-\frac{2a_2}{a_4}}$$

$a_4 > 0 \rightarrow$ physical

only Gou fixed pm soln

is accepted.

W.F. is rejected

For $d < 4$ ($\epsilon > 0$)

$$\text{W.F. } (r^*, u_{\frac{x}{4}}^*) = \left(-\frac{|\epsilon|\Lambda^2}{6}, +\frac{2\pi^2|\epsilon|}{9} \right)$$

$$u^* = \frac{2\pi^2|\epsilon|}{9} > 0 \rightarrow a_4 > 0$$

Physical solution

at W.F. $y_2 = -\epsilon \rightarrow$ Irrelevant.
 $d < 0$

At W.F. $\epsilon = 4-d$

$$\left\{ \begin{array}{l} \Lambda_t = l^{(2-\epsilon/3)} = l^{2(1-\epsilon/6)} \rightarrow y_t = 2(1-\epsilon/6) \\ \Lambda_2 = l^{-\epsilon} \Rightarrow y_2 = -\epsilon \end{array} \right.$$

For $d > 4$ ($\epsilon < 0$)

$$y_t = 2 + \frac{|\epsilon|}{3} > 2 \rightarrow \text{Relevant.}$$

$y_2 = +|\epsilon| \rightarrow$ Relevant *

$$\text{W.F. } \left(-\frac{\epsilon\Lambda^2}{6}, \frac{2\pi^2\epsilon}{9} \right) = \left(\frac{+|\epsilon|\Lambda^2}{6}, -\frac{2\pi^2|\epsilon|}{9} \right)$$

No physical solution

For $d < 4$ ($\epsilon > 0$) $\gamma_t = 2 - \frac{|\epsilon|}{3}$

$\gamma_2 = -|\epsilon| < 0 \rightarrow$ Irrelevant attractive Fixed point

W.F. $(-\frac{|\epsilon|\Lambda^2}{6}, +\frac{2\pi^2|\epsilon|}{9})$ physical solution

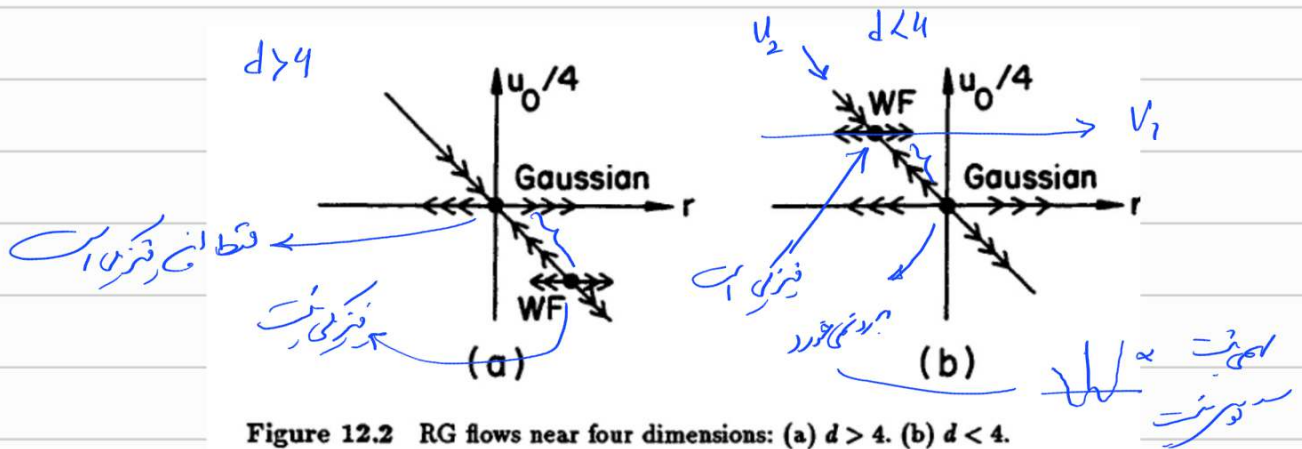


Figure 12.2 RG flows near four dimensions: (a) $d > 4$. (b) $d < 4$.

Notice that

$$\left(r_{WF}^*, \frac{u_{WF}^*}{4} \right) = \left(-\frac{\epsilon \Lambda^2}{6}, \frac{2\pi^2 \epsilon}{9} \right), \quad V_2 = \left(-\frac{3\Lambda^2}{4\pi^2}, 1 \right)$$

$$= \frac{2\pi^2 \epsilon}{9} V_2 = \frac{2\pi \epsilon}{9} \left(-\frac{3\Lambda^2}{4\pi^2}, 1 \right)$$

For W.F. \rightarrow Gaussian fixed-point

$\epsilon \rightarrow 0$

For $d > 4$ ($\epsilon < 0$) W.F. has negative value of quartic coupling constant which is unphysical.

For $d < 4$ ($\epsilon > 0$)

W. F. has positive value of
quartic coupling const
which is physical.

⑧
$$\nu = \frac{1}{\gamma_t} = \frac{1}{2} \left(1 + \frac{\epsilon}{6} \right) + O(\epsilon^2)$$

$\gamma = 1 + \epsilon/6$

$\eta = 0$

$\alpha = \epsilon/6$

$\beta = \frac{1}{2} - \epsilon/6$

$\delta = 3 + \epsilon$

} $O(\epsilon^2)$

Table 12.1 CRITICAL EXPONENTS FOR THE ISING UNIVERSALITY CLASS IN THREE DIMENSIONS

Exponent	ϵ -expansion to $O(\epsilon)$	Mean Field	Experiment	Ising ($d = 3$)
α	0.167	0 (disc.)	0.110 - 0.116	0.110(5)
β	0.333	1/2	0.316 - 0.327	0.325 ± 0.0015
γ	1.167	1	1.23 - 1.25	1.2405 ± 0.0015
δ	4.0	3	4.6 - 4.9	4.82(4)
ν	0.583	1/2	0.625 ± 0.010	0.630(2)
η	0	0	0.016 - 0.06	0.032 ± 0.003

ϵ -Expansion

Meanfield

$d < 4 \rightarrow$ Ginzburg criterion

shows that $d > 4$

for mean field

For α

	$n=1$	$n=2$	$n=3$	$n=4$
$\mathcal{O}(\epsilon)$ at $\epsilon=1$	0.17	0.11	0.06	0
Experiments in $d=3$	<u>0.11</u>	-0.01	-0.12	-

From Page 90
Kardar

$$\nu = \frac{1}{\gamma_k} = \frac{1}{2} + \frac{1}{4} \frac{n+2}{n+8} \epsilon + \mathcal{O}(\epsilon^2) \quad (5.60)$$

$$\alpha = 2 - d\nu = \frac{4-n}{2(n+8)} \epsilon + \mathcal{O}(\epsilon^2) \quad (5.61)$$

$$\beta = \frac{1}{2} - \frac{3}{2(n+8)} \epsilon + \mathcal{O}(\epsilon^2) \quad (5.63)$$

$$r = \frac{2\gamma_{n-d}}{\gamma_k} = 1 + \frac{n+2}{2(n+8)} \epsilon + \mathcal{O}(\epsilon^2) \quad (5.64)$$



برای آن ادبیات مطابق مطابق؛ تحلیل ابعاد مطابق؛
Sec 4.7
Kardar

Eq. (4.64)

$$Z_s = Z_0 Z_m$$

Coarse-grain ϵ \uparrow
Regular

$\left. \begin{array}{l} \text{rescale} \\ \text{renormalize} \end{array} \right\} \quad (4.67)$

$$\beta_{\text{eff}} = \left[\int_0^{\Lambda \ell} \frac{d^d q}{(2\pi)^d} \left(\frac{t + Kq^2 + Lq^4 + \dots}{2} \right) |m(q)|^2 \right] - h \cdot \bar{m}(q=0)$$

↓

$$\left. \begin{array}{l} \text{Rescaling} \rightarrow g' = g \ell \\ \text{Renormalization} \rightarrow m' = m/\ell \end{array} \right\} \rightarrow (4.67)$$

