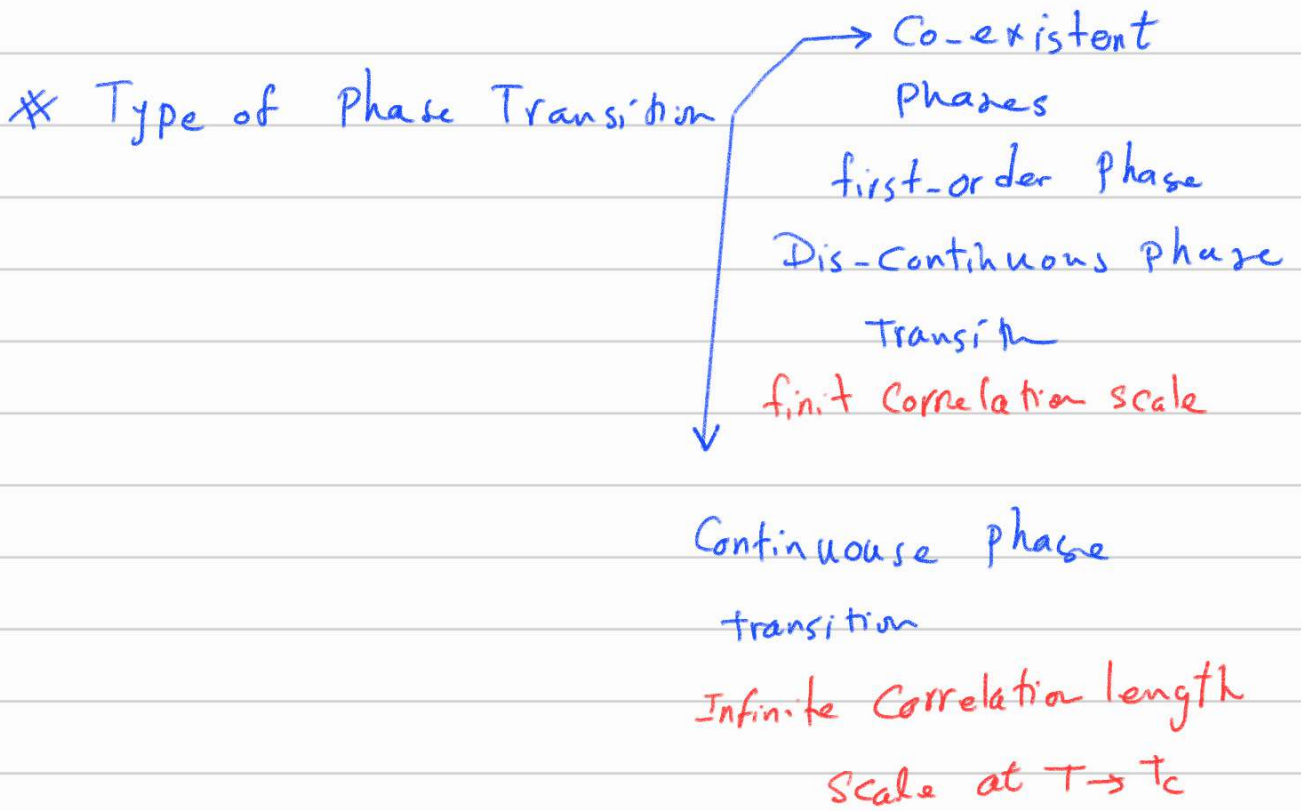


Chapter 2: Phase Transition in Simple System (Cardy)



§ : Correlation Length Scale
Time
Angular Scale.

Correlation function

تابع همبستگی

{

 → ارجاع می دهیم در این فرآیندها که صادر می شود

 Weigthed Two Point Correlation function

 وزن دار دو نقطه ای تابع خود همبستگی

 Auto-Correlat

 }

$$C_{\delta S}(\vec{r}_i, \vec{r}_j) = \langle \delta S(\vec{r}_i) \delta S(\vec{r}_j) \rangle$$

$$\delta S = G(\vec{r}_i, \vec{r}_j) \rightarrow \text{Connected moment}$$

$M = \text{Magnetization}$

$$\delta S(\vec{r}_i) = S(\vec{r}_i) - \langle S(\vec{r}_i) \rangle_r$$

$$G(\vec{r}_i, \vec{r}_i) = \sigma_0^2$$

Homogeneity system

$$G(\vec{r}_i, \vec{r}_j) = G(r)$$

$$r = |\vec{r}_i - \vec{r}_j|$$

$G(\vec{r}_i, \vec{r}_j)$: Response function

\vec{E}

Exercise 1:

$$\text{for } \mathcal{H} = \mathcal{H}(\text{fst}) - \sum_{i=1}^N H_i S_i$$

Ideal Gas

$$\mathcal{H} = \sum \frac{p_i^2}{2m} + \sum_{ij} U_{ij}$$

Magnetic field
at S_i

$$G(r_i, r_j) = \langle \delta S(i) \delta S(j) \rangle$$

$$= K_B T \frac{\partial}{\partial H_i} \langle S_j \rangle$$

تغییرات در محل S_j به دلیل تغییرات در محل r_i

Susceptibility χ : Global Response function

$$\sum_{ij} G(\vec{r}_i, \vec{r}_j) \propto \chi$$

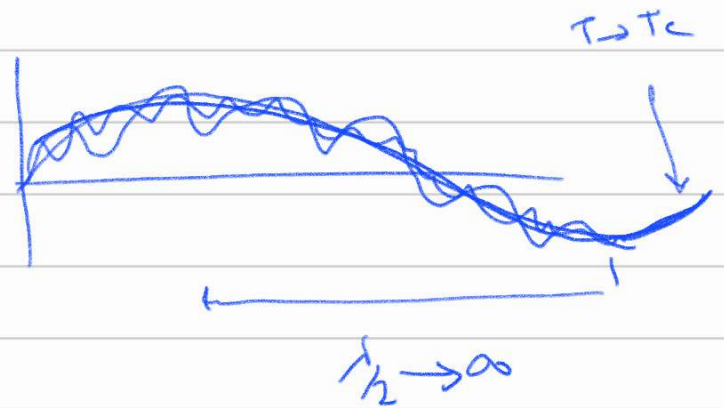
$T \rightarrow T_c$ χ : Static Susceptibility

Dominant Mode for variation of spins

$$\chi(\omega \rightarrow 0) \quad (K \rightarrow \infty) \quad T \rightarrow \infty$$

توضیح

Langern Model building



Effective Dynamical model.

$$\left\{ \frac{d\langle \vec{S}_i \rangle}{dt} = ? \leftarrow \text{Langern} \right\}$$

$T \rightarrow T_c$

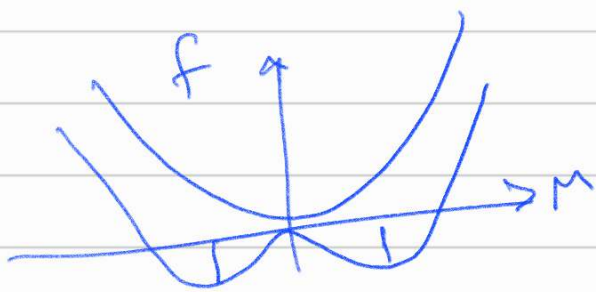
توضیح

Free Energy \rightarrow

$$f = a + bT M^2 + cM + dHM + \dots + Q(M^6)$$

$$\frac{\partial f}{\partial M} \Big|_{T=T_c} = 0$$

$$T > T_c \rightarrow M \left[F = -k_B T \ln Z \right]$$



$$G(r) = \boxed{r = |\vec{r}_i - \vec{r}_j|}$$

$$\chi \propto \int d^d r G(r)$$

$$\boxed{\chi \propto \xi^2}$$

If $d=3$ $G(r) = \frac{1}{r} f(r/\xi)$

$\eta=0$

$$\chi \propto \xi^{2-\eta} \text{ (universal)}$$

Unit cell character

$$G(r) = \frac{1}{r} f(r/\xi, a/\xi)$$

$$f(r/\xi, a/\xi) \sim (a/\xi)^2$$

$$\boxed{\chi \propto a^\eta \xi^{2-\eta}}$$

Mean field Theory $\underline{\underline{\eta=0}}$

General Properties of Ising Model.

$$\mathcal{H} = -H \cdot \sum_{i=1}^N \vec{S}_i - \sum_{\langle ij \rangle} \bar{J}_{ij} \vec{S}_i \cdot \vec{S}_j$$



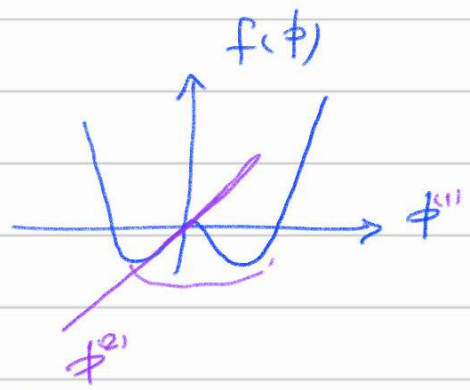
$$M = \frac{1}{N} \sum_{i=1}^N \langle S_i \rangle$$

$$i: \phi_i = \{\phi_i^1, \phi_i^2, \dots, \phi_i^n\}$$

$$M = -\frac{1}{N} \partial \frac{f}{\partial H} \quad \text{--- } O(n)$$

$$O(2) = \{ \phi^{(1)}, \phi^{(2)} \}$$

$$O(1) = \{ \phi \}$$



$$O(1) \Rightarrow f = a + bt\phi^2 + c\phi^4 + \dots$$

$$O(n) \Rightarrow f = a + bt \sum_{a=1}^n (\phi^{(a)})^2 + c \sum_{a=1}^n (\phi^{(a)})^4 + \dots$$

free Energy.

(A) $f < 0$

(B) $f(H, T, J, \dots)$ is continuous

(C) $\frac{\partial f}{\partial T}, \frac{\partial f}{\partial H}, \dots$ exist everywhere

(D) $S'(\text{Entropy}) = - \frac{\partial f}{\partial T} \geq 0$

(E) $\frac{\partial f}{\partial T}$ is monotonically non-increasing.

$$\frac{\partial^2 f}{\partial T^2} \leq 0$$

$$C \propto - \frac{\partial^2 f}{\partial T^2} \geq 0$$

$$\frac{\partial f}{\partial H} \dots \rightarrow \frac{\partial^2 f}{\partial H^2} \leq 0$$

$$\chi \propto - \frac{\partial^2 f}{\partial H^2} \geq 0$$

$$f \propto \leftarrow \rightarrow$$

$$f = \frac{F_{\Omega}}{\bar{V}(\Omega)} = -k_B T \ln Z_{\Omega}$$

$$Z_{\Omega} = \sum_{\{s\}} e^{-\beta H} = 0.01 + 0.001 + 1 + \dots$$

صافتی یک سیرینده وجود دارد که همین است که کم است یعنی $H < 0$

$$e^{-\beta H} > 1 \rightarrow Z > 1 \rightarrow F_{\Omega} < 0$$

Recall scaling Exponents for
Ising model.

{	$M \propto t ^{-\beta}$ for $H \rightarrow 0$	β γ δ α ν η
	$\chi \propto t^{-\gamma}$	
	$M \propto H^{1/\delta}$ $t \rightarrow 0$	
	$C \propto t ^{-\alpha}$	
	$\xi \propto t ^{-\nu}$	
$G(r) \propto \frac{1}{r^{d-2+\eta}}$		

fluid : A classical fluid

Including N -Particles
at $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$

$$H = H_{\text{Kinetic}} + H_{\text{Interaction}} \quad \text{--- (int)}$$

$$= \dots + \sum_{i,j} U(\vec{r}_i - \vec{r}_j) = \dots - \sum_{i,j} J_{ij} n(r_i) n(r_j)$$

$$Z(T, V, z) \propto \sum_{N=0}^{\infty} \frac{z^N}{N!} \int d^3r_1 d^3r_2 \dots d^3r_N e^{-\beta H_{\text{Int.}}}$$

fugacity $z = e^{\beta \mu}$

Mapping \rightarrow Lattice Gas

Occupation No. $n(r)$ رض کنيد قضا، البته مي بينم . عدد اشغال

$$n(r_i) = \begin{cases} 1 \\ 0 \end{cases}$$

Constraint $\sum_r n(r) = N$

$$Z(T, V, z) = \sum_{N=0}^{\infty} \sum_{\{n(r)\}} \frac{1}{N!} \dots$$

$$= \sum_{\{n(r)=0,1\}} e^{\beta \mu \sum_r n(r)} \frac{1}{e^{-\beta(-2) \sum_{rr'} J_{rr'} n(r) n(r')}} z^N$$

$$\rightarrow S(r) \equiv 2n(r) - 1 \rightarrow S(r) \begin{cases} -1 \\ +1 \end{cases}$$

$$Z = \sum_{\{n(r)\}} e^{-2\beta \sum_{r,r'} J_{rr'} n(r)n(r')}$$

$$= Z^{N/2} Z^{\sum \frac{S(r)}{2}} e^{-2\beta \sum_{r,r'} J_{rr'} \left(\frac{S(r)+1}{2}\right) \left(\frac{S(r')+1}{2}\right)}$$

$$Z = \sum_{\{S(r) = -1, +1\}} e^{\frac{\beta}{2} \sum_{r,r'} J_{rr'} S(r)S(r') + \beta H \sum S(r)}$$

$$H \equiv \sum_{r'} J_{rr'} + \frac{K_B T}{2} \ln Z$$

$N(r)=0$ ~~از این محاسبه~~
 $N(r)=+1$ ~~از این محاسبه~~
 معادله یک مدل آنتروپی، یک مدل خارجی در نقطه r

$$H(r)_{\text{ext}} = \frac{1}{2} K_B T \ln Z + \sum_{r'} J_{rr'} S(r')$$

\uparrow
 $e^{\beta \mu}$

Mean field Theory

Chapter 2: Cardy

